

A numerical introduction to tensor networks for quantum simulation

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[Credit for course materials: Prof. Jan von Delft]

1. Why tensor networks?

Tensor networks provide a flexible description of quantum states.

In some cases, they are efficient - can accurately describe state with polynomial resources.

Example: spin- chain, with sites:

Local state space of site :

Local state label:

Local dimension:

Shorthand:

Index on state label suffices to identify the site Hilbert space

Generic basis state for full chain of length N
(convention: add state spaces for new sites from left):

Hilbert space for full chain:

Generic quantum state:

Dimension of full Hilbert space : (# of different configurations of)

Specifying involves specifying , i.e. different complex numbers.

is a tensor of rank (rank = # of legs)

Graphical representation: (one leg for each index)

Claim: such a rank L tensor can be represented in many different ways:

MPS: matrix product state

PEPS: projected entangled-pair state

arbitrary tensor network

- a link between two sites represents entanglement between them
- different representations \rightarrow different entanglement book-keeping
- tensor network = entanglement representation of a quantum state

2. Iterative diagonalization

Consider a spin- s chain with Hamiltonian

local state space for site :

We seek eigenstates of :

Diagonalize Hamiltonian iteratively, adding one site at a time.

$N = 1$: Start with first site, diagonalize H_1 in Hilbert space \mathcal{H}_1 . Eigenstates have form

(sum over i implied) $H_1 = \sum_i \epsilon_i |i\rangle\langle i|$ coefficient matrix $H_{ij} = \epsilon_i \delta_{ij}$ combining 'incoming' $|i\rangle$ into

$N = 2$: Add second site, diagonalize H_2 in Hilbert space \mathcal{H}_2 :

(sum over i, j implied) $H_2 = \sum_{i,j} H_{ij} |i\rangle\langle j|$ coefficient matrix $H_{ij} = \epsilon_i \delta_{ij} + t \delta_{i,j\pm 1}$ combining 'incoming' $|i\rangle$ into

$N = 3$: Add third site, diagonalize H_3 in Hilbert space \mathcal{H}_3 :

Your try:

Answer:

Continue similarly until having added site N. Eigenstates of structure: have following

Nomenclature: = physical indices, = (virtual) bond
indices

Alternative, widely-used notation: 'reshape' coefficient tensors as

to highlight 'matrix product' structure in noncovariant notation:

Comments

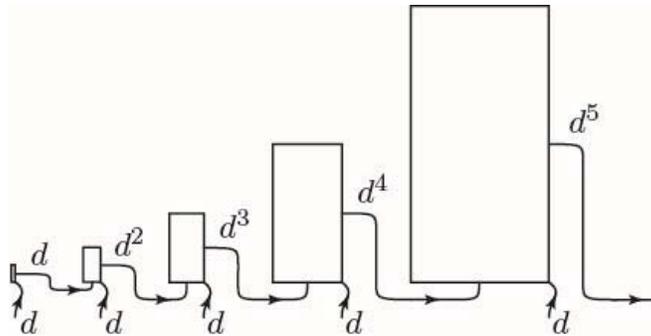
1. Iterative diagonalization of 1D chain generates eigenstates whose wave functions are tensors that are expressed as matrix products -> matrix product states (MPS)

Matrix size grows exponentially.

for given d , d^2 has dimension d^2 (vector)

for given d , d^3 has dimension d^3 (rectangular matrix)

for given d , d^5 has dimension d^5 (larger rectangular matrix)



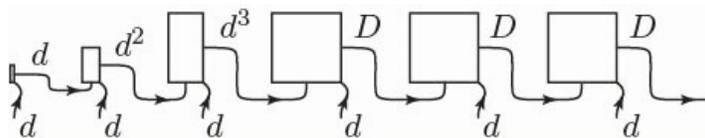
“Hilbert space is a large place!”

Numerical costs explode with increasing N , so truncation schemes are needed.

Truncation can be done in a controlled way using tensor network methods.

Standard truncation scheme: use

for all virtual bonds



2. Number of parameters available to encode state:

[equal if all virtual bonds have same dimension D]

If N is large:

Why should we expect this ansatz with polynomial parameters to accurately represent a quantum state?

Remarkable fact: for 1d Hamiltonian with local interactions and a gapped spectrum, its ground state can be accurately represented by MPS!

Reason: Area law for entanglement entropy. We will discuss later.

3. Covariant index notation

More detail of covariant index notation is in L2 & L10 of “Mathematics for Physicists”, Altland & von Delft, see [here](#).

Kets (Hilbert space vectors)

For kets, indices are down. E.g. basis kets:

For components of kets (wrt a basis), indices are upper:

Repeated indices (always up-down pairs) are always summed (implied summation).

Example: linear combination of kets.

Note: for $\langle i |$ the index i identifies components of kets \rightarrow upper
 for $| j \rangle$ the index j identifies components of basis kets \rightarrow lower

Basis for direct product space:

Note ket order: start with first space on very right, successively attach new spaces from left.

Linear combinations:

Bras (Hilbert space dual vectors)

For bras, indices are upper. E.g. basis bras:

For components of bras (wrt a basis), indices are lower:

Complex conjugation:

Linear combinations of bras:

Complex conjugation:

Note: for $\langle i |$, the index i identifies basis bras (dual vectors) hence upper
 lower $| j \rangle$ the index j identifies components of bras (dual vectors) hence

Basis for direct product space:

Note bra order: opposite to kets so expectation values yield nested bra-ket pairs:

Linear combinations:

Complex conjugation:

Orthonormality

If $\{ | \psi_i \rangle \}$ form orthonormal basis:

If $\{ \langle \psi_i | \}$ form orthonormal basis, too:

Combined:

Hence A is unitary:

Operators

Simplified notation

It is customary to simplify notational conventions for kets and bras.

In kets, use subscript indices as ket names:

In bras, use superscript indices as bra names:

Now up/down convention for indices is no longer displayed but it still implicit!

Linear combination of kets:

Coefficient matrix = overlap:

If direct products are involved:

Coefficient matrix = overlap:

Linear combination of bras:

Coefficient matrix = overlap:

If direct products are involved:

Coefficient matrix = overlap:

Operators:

In the overlaps:

bra indices: upper on or , as incoming arrows

ket indices: lower on or , as outgoing arrows

4. Entanglement entropy and Area Laws (introductory comments)

Consider a quantum system in state $|\psi\rangle$, with density matrix ρ

Divide system into two parts, A and B . Suppose A has linear dimension l .

To obtain reduced density matrix of A (or B), trace out B (or A)

'reduced density matrix' for A :

'Entanglement entropy' of A and B :

Key result: for Hamiltonians with only local interactions, ground state is governed by an 'area law':

(area of boundary of)

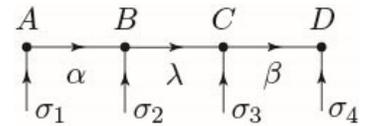
in 3D for gapped system

in 2D for gapped system

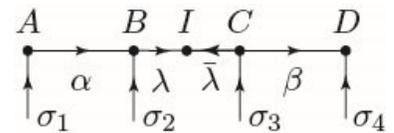
in 1D for gapped system

in 1D for *gapless system

Now consider an MPS of maximal bond dimension D :



Divide system into two parts: Left -> 2 sites, Right -> 2 sites



=entangled superposition of two state spaces,
each having dimension of at most D

(After the sum over α has been performed explicitly using
Kronecker delta, result contains non-covariantly paired indices

Density matrix:

Reduced density matrix:

with matrix elements

This matrix has rank D (say D)

Let λ_i be its eigenvalues, with

and normalization

Entanglement entropy:

Maximal if $S_A = \log d_B$ for all A :

1D gapped:

1D critical:

2D gapped:

3D gapped:

Conclusion: MPS can encode ground state efficiently for gapped and gapless systems in 1D, but not 2D or 3D!

5. Tensor network diagrams

[Orus 2014, Sec 4.1]

'tensor' = multi-dimensional array of numbers

'rank of tensor' = number of indices = # of legs

rank-0: scalar

rank-1: vector

rank-2: matrix

rank-3: tensor

Index contraction: summation over repeated index

$$\begin{array}{c} \xrightarrow{\alpha} \bullet \xrightarrow{\gamma} \\ C \end{array} = \begin{array}{c} \xrightarrow{\alpha} \bullet \xrightarrow{\beta} \bullet \xrightarrow{\gamma} \\ A \quad B \end{array}$$

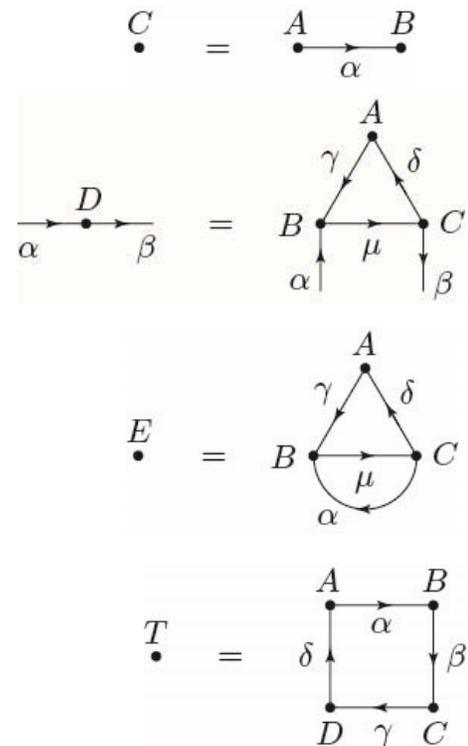
graphical representation of matrix product

= 'bond dimension' of index

'open index' = non-contracted index (here)

'tensor network' = set of tensors with some or all indices contracted according to some pattern

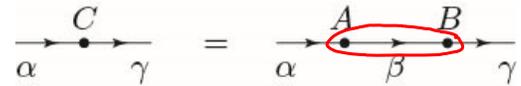
Examples:



Cost of computing contractions

Result of contraction does not depend on order in which indices are summed, but numerical cost does!

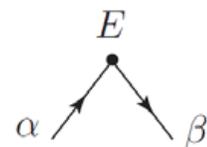
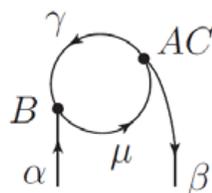
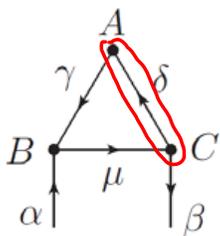
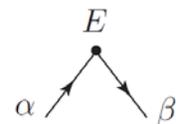
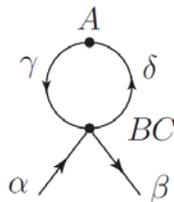
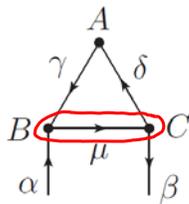
Example 1: cost of matrix multiplication is



For every fixed values of α and γ (combinations), sum over β

Cost = $\sum_{\beta} C_{\alpha\beta\gamma}$ (simplifies to $\alpha\gamma$ if all bond dimensions are equal)

Example 2:



First contraction scheme has total cost _____, second has _____ !

Finding optimal contraction order is difficult. In practice, trial and error....

In first part of course we will focus on 1D tensor networks, then 2D.

6. Singular value decomposition (SVD)

[Schollwoeck2011, Sec 4]

Any matrix _____ of dimension _____ can be written as

$$\begin{aligned}
 D \leq D': \quad & \begin{matrix} D' \\ D \end{matrix} \begin{matrix} \text{[Grey Box]} \end{matrix} = \begin{matrix} D & D & D' \\ D \begin{matrix} \text{[Vertical Lines]} \end{matrix} & \begin{matrix} \text{[Diagonal Box]} \end{matrix} & \begin{matrix} \text{[Horizontal Lines]} \end{matrix} \\
 & M = U \quad S \quad V^\dagger \\
 D \geq D': \quad & \begin{matrix} D' \\ D \end{matrix} \begin{matrix} \text{[Grey Box]} \end{matrix} = \begin{matrix} D' & D' & D' \\ D \begin{matrix} \text{[Vertical Lines]} \end{matrix} & \begin{matrix} \text{[Diagonal Box]} \end{matrix} & \begin{matrix} \text{[Horizontal Lines]} \end{matrix}
 \end{aligned}$$

Properties of S

- square matrix of dimension _____
- diagonal with non-negative diagonal elements (singular values)
- Schmidt rank _____ : number of non-zero singular values

- arrange in descending order:

Properties of _____:

- matrix of dimension

$$D \leq D': \quad \begin{array}{c} U^\dagger \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \begin{array}{c} U \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$D_{\min} = D$

- columns are orthonormal

$$D \geq D': \quad \begin{array}{c} U^\dagger \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \begin{array}{c} U \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

$D_{\min} = D'$

Properties of _____:

- matrix of dimension

$$D \leq D': \quad \begin{array}{c} V^\dagger \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \begin{array}{c} V \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

- rows are orthonormal

$$D \geq D': \quad \begin{array}{c} V^\dagger \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \begin{array}{c} V \\ \text{---} \\ \text{---} \\ \text{---} \end{array} = \begin{array}{c} 1 \\ \text{---} \\ \text{---} \\ \text{---} \end{array}$$

Useful properties:

So columns of U are eigenvectors of $A^T A$

Columns of V are eigenvectors of $A A^T$

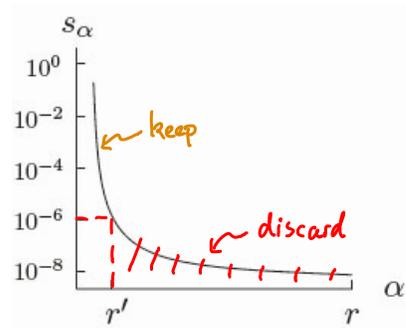
Truncation

SVD yields optimal approximation of rank r matrix A by a rank r matrix A_r

(optimal wrt Frobenius norm)

Suppose

with



QR decomposition

If singular values are not needed,

a matrix

has the 'full QR decomposition'

$$\begin{matrix}
 & D' \\
 D & \boxed{} \\
 \\
 & M \\
 \\
 & D' \\
 D & \boxed{}
 \end{matrix}
 =
 \begin{matrix}
 & D & D' \\
 D & \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix} & \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix} \\
 \\
 & Q & R \\
 \\
 & D & D' \\
 D & \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix} & \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}
 \end{matrix}$$

with a unitary matrix

and a upper triangular matrix

If $D \geq D'$, then M has the 'thin QR decomposition'

$$\begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}
 \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}
 \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}
 =
 \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}
 \begin{matrix} \boxed{} \\ \boxed{} \\ \boxed{} \\ \boxed{} \end{matrix}$$

with $\dim(Q_1) =$, $\dim(R_1) =$,

and R1 upper triangular.

QR is numerically cheaper than SVD but has less information (does not provide rank).